

# CONVERGENCE OF THE EDIIS ALGORITHM FOR NONLINEAR EQUATIONS

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**Abstract.** The EDIIS (Energy Direct Iteration on the Iterative Subspace) algorithm was designed to globalize Anderson acceleration, a method for improving the performance of fixed point iteration. The motivating application is electronic structure computations. In this paper we prove a convergence result for that algorithm and illustrate the theory with a computational example.

**Key words.** Nonlinear equations, Anderson acceleration, EDIIS

**AMS subject classifications.** 65H10, 81V55

**1. Introduction.** The purpose of this paper is to analyze the convergence of the EDIIS (Energy Direct Inversion on the Iterative Subspace) algorithm [21]. EDIIS is a modification of Anderson acceleration [1] or the DIIS (Direct Inversion on the Iterative Subspace) method [21, 22, 34, 37]. EDIIS relaxes the need for a sufficiently accurate initial iterate. EDIIS is the default solver for the SCF (self consistent field) iteration in the widely-used Gaussian [12] quantum chemistry code. We prove convergence from any starting point in a convex set in which the fixed point map is a contraction and then analyze local convergence. Our local convergence is an improvement of the result in [41] and applies to both EDIIS and Anderson acceleration.

We will begin this introductory section with a review of Anderson acceleration and some of the recent convergence results. We will then describe the EDIIS algorithm. In § 2 we prove our convergence results. Finally, in § 3 we will report on a computation which both illustrates the theory and, as is also done in [21], shows how the convergence speeds for EDIIS and Anderson acceleration, while identical in theory, can differ significantly in practice.

Our notational convention is that vectors and vector-valued functions in  $R^N$  are in bold. Scalars and elements of infinite dimensional spaces (eg integral operators and the functions acted upon by those operators) are in the usual italic math font.

Anderson acceleration [1] is an iterative method for fixed point problems of the form

$$(1.1) \quad \mathbf{u} = \mathbf{G}(\mathbf{u}),$$

where  $\mathbf{u} \in R^N$  and  $\mathbf{G} : R^N \rightarrow R^N$ . The method was designed to accelerate Picard or fixed point iteration *i.e.*

$$(1.2) \quad \mathbf{u}_{k+1} = \mathbf{G}(\mathbf{u}_k).$$

Anderson acceleration was originally designed for integral equations and has been widely used in electronic structure computations (see [9] and many references since then) and is now very common in that field. Anderson acceleration is essentially the same as Pulay mixing [32, 33], DIIS [21, 22, 34, 37], nonlinear GMRES [4, 25, 30, 45]. Other applications include nuclear reactor design [16, 42], stiff dynamics [13], hydrology [24], and fluid-structure interaction [10, 15, 23] where the method is called interface quasi-Newton.

The analysis of Anderson acceleration is far from complete. In this paper we assume, as do all theoretical results about this algorithm, that the map  $\mathbf{G}$  is a contraction. In practice, however, Anderson acceleration does very well for problems in which  $\mathbf{G}$  is either definitely not a contraction [41] or not provably a contraction. The results here do not explain those cases.

Anderson acceleration was designed for a problem where Newton's method is not practical because obtaining approximate Jacobians or Jacobian-vector products is too costly. One should expect that Newton's

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method would perform better when derivative information can be had at reasonable cost and we have certainly found that to be the case in our own recent work [16]. Anderson iteration maintains a history of residuals

$$\mathbf{F}(\mathbf{u}) = \mathbf{G}(\mathbf{u}) - \mathbf{u}$$

of size at most  $m + 1$ , where the **depth**  $m$  is an algorithmic parameter. When  $m$  is important, we will call the iteration  $\text{Anderson}(m)$ .  $\text{Anderson}(0)$  is Picard iteration by definition.

The formal description in Algorithm 1 is most convenient for analysis and exposition, but not for implementation. We refer to [7, 38, 39, 41, 43, 44] for examples of efficient implementations.

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**Algorithm 1** Anderson Acceleration

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$\text{anderson}(\mathbf{u}_0, \mathbf{G}, m)$

$\mathbf{u}_1 = \mathbf{G}(\mathbf{u}_0); \mathbf{F}_0 = \mathbf{G}(\mathbf{u}_0) - \mathbf{u}_0$

**for**  $k = 1, \dots$  **do**

    Choose  $m_k \leq \min(m, k)$

$\mathbf{F}_k = \mathbf{G}(\mathbf{u}_k) - \mathbf{u}_k$

    Minimize  $\left\| \sum_{j=0}^{m_k} \alpha_j^k \mathbf{F}_{k-m_k+j} \right\|$  subject to  $\sum_{j=0}^{m_k} \alpha_j^k = 1$

$\mathbf{u}_{k+1} = \sum_{j=0}^{m_k} \alpha_j^k \mathbf{G}(\mathbf{u}_{k-m_k+j})$

**end for**

---

The iteration uses the most recent  $m + 1$  residuals  $\mathbf{F}(\mathbf{u}_j)$  for  $k - m_k \leq j \leq k$  where  $m_k \leq \min(k, m)$ . The key step in the iteration is solving the **optimization problem**

$$(1.3) \quad \text{Minimize} \left\| \sum_{j=0}^{m_k} \alpha_j^k \mathbf{F}(\mathbf{u}_{k-m_k+j}) \right\| \text{ subject to } \sum_{j=0}^{m_k} \alpha_j^k = 1,$$

for the coefficients  $\{\alpha_j^k\}$ .

Any vector norm can be used in the optimization problem with no change in the convergence theory [41]. In particular the optimization problem for the coefficients in either the  $\ell^1$  or  $\ell^\infty$  norms can be formulated as a linear programming problem [8]. The optimization problem is easier to solve if one uses the  $\ell^2$  norm and that is standard practice. In this case optimization problem for the coefficients can be expressed as a linear least squares problem and solved very inexpensively. One way to do this is to solve the linear least squares problem

$$(1.4) \quad \text{Minimize} \left\| \mathbf{F}(\mathbf{u}_k) + \sum_{j=0}^{m_k-1} \alpha_j^k (\mathbf{F}(\mathbf{u}_{k-m_k+j}) - \mathbf{F}(\mathbf{u}_k)) \right\|_2^2,$$

for  $\{\alpha_j^k\}_{j=0}^{m_k-1}$ . Then one recovers  $\alpha_{m_k}^k$  by

$$\alpha_{m_k}^k = 1 - \sum_{j=0}^{m_k-1} \alpha_j^k.$$

The choice of  $m_k$  is, in the original form, simply  $\min(m, k)$ . One can adapt  $m_k$  as the iteration progresses to, for example, enforce well-conditioning of the linear least squares problem (1.4) [39, 44].

One can also [11, 31, 34, 35, 44] show that Anderson acceleration is related to multiseccant quasi-Newton methods or, in the case of linear problems, GMRES. None of these results lead to a convergence proof, even in the linear case, unless the available storage is large enough to allow GMRES to take a number of iterations equal to the dimension of the problem. The recent work of one of the authors and his students [39–41] contains the first convergence theory for Anderson acceleration as it is applied in practice.

**1.1. Convergence Theory.** Theorem 1.1 is one of the convergence results from [41]. That paper also has results for several special cases. We assume that  $\mathbf{G}$  is a contraction with contractivity constant  $c \in (0, 1)$  in a closed set  $D \subset R^N$ ,

$$(1.5) \quad \|\mathbf{G}(\mathbf{u}) - \mathbf{G}(\mathbf{v})\| \leq c \|\mathbf{u} - \mathbf{v}\|$$

72 for all  $\mathbf{u}, \mathbf{v} \in D$ . The contraction mapping theorem implies that  $\mathbf{G}$  has a unique fixed point  $\mathbf{u}^* \in D$ . As is  
 73 standard, we let  $\mathbf{e} = \mathbf{u} - \mathbf{u}^*$  and make the assumption from [41] on the smoothness of  $\mathbf{G}$  and the Anderson  
 74 iteration coefficients.

75 The convergence of the Picard iteration for a contraction map is q-linear [19] with q-factor  $c$  *i. e.*

76 
$$\|\mathbf{e}_k\| \leq c\|\mathbf{e}_{k-1}\|.$$

77 We will show in this paper that Anderson acceleration is r-linear with r-factor  $c$ , which means

78 
$$\|\mathbf{e}_k\| = O(c^k).$$

79 ASSUMPTION 1.1.  $\mathbf{G}$  is a Lipschitz continuously differentiable in the ball

80 
$$\mathcal{B}(\hat{r}) = \{\mathbf{u} \mid \|\mathbf{e}\| \leq \hat{r}\} \subset D,$$

81 for some  $\hat{r} > 0$ .

82 There is  $M_\alpha$  such that for all  $k \geq 0$

83 
$$\sum_{j=0}^{m_k} |\alpha_j^k| \leq M_\alpha.$$

84 The differentiability assumption is needed in the analysis, but not in the formulation or implementation  
 85 of the algorithm. Our convergence result in § 2.2 relaxes the assumption to continuous differentiability.

86 THEOREM 1.1. [41] Let Assumption 1.1 hold and let  $c < \hat{r} < 1$ . Then if  $\mathbf{u}_0$  is sufficiently close to  $\mathbf{u}^*$ ,  
 87 the Anderson iteration converges to  $\mathbf{u}^*$ . In fact, for all  $k \geq 0$ ,

88 (1.6) 
$$\|\mathbf{F}(\mathbf{u}_k)\| \leq \hat{r}^k \|\mathbf{F}(\mathbf{u}_0)\| \text{ and } \|\mathbf{e}_k\| \leq \left(\frac{1+c}{1-c}\right) \hat{r}^k \|\mathbf{e}_0\|.$$

89 The interpretation of this result is that if the initial data are sufficiently good, then the r-factor for Anderson  
 90 iteration is no worse than the q-factor of Picard iteration as predicted by the contractivity constant  $c$ . While  
 91 r-linear convergence is weaker than q-linear, Anderson acceleration is often faster than Picard iteration in  
 92 practice. The requirement that the initial iterate be near the solution is also meaningful in practice [36, 46, 47]  
 93 and motivated the EDIIS algorithm [21] which is the subject of this paper.

94 Both Picard iteration and Anderson acceleration can perform better than the prediction (see § 3). In  
 95 practice, Anderson acceleration is often significantly better than Picard iteration, but there is no theory that  
 96 explains this under practical (*i. e.* very limited storage) operating conditions.

97 **1.2. The EDIIS Algorithm.** Anderson acceleration performs poorly for some applications. One  
 98 example is electronic structure computations for metallic systems where the HOMO-LUMO gap is small  
 99 and a good initial iterate is difficult to obtain. In this case both Picard iteration and Anderson acceleration  
 100 perform poorly [21]. In such cases one can sometimes use a small mixing parameter to ensure convergence,  
 101 especially when the initial iterate is poor. However, a small mixing parameter may degrade the performance  
 102 of the iteration especially when near the solution. The role of the damping parameter  $\beta$  in Picard iteration  
 103 is simple damping

104 
$$\mathbf{u}_{k+1} = (1 - \beta)\mathbf{u}_k + \beta\mathbf{G}(\mathbf{u}_k).$$

105 If one applies EDIIS or Anderson acceleration to

106 
$$\mathbf{G}_\beta(\mathbf{u}) = (1 - \beta)\mathbf{u} + \beta\mathbf{G}(\mathbf{u}).$$

107 then [40] one obtains

108 
$$\mathbf{u}_{k+1} = (1 - \beta) \sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j} + \beta \sum_{j=0}^{m_k} \alpha_j^k \mathbf{G}(\mathbf{u}_{k-m_k+j}),$$

109 which is how damping is done in Anderson acceleration [1].

110 One attempt to solve these problems for small systems is the EDIIS algorithm from [21]. In [21] the  
 111 authors also formulated the fixed point problem to directly minimize energy, hence the name of the method,  
 112 but that does not affect the convergence analysis in this paper.

113 EDIIS differs from Anderson acceleration by imposing a nonnegativity constraint on the coefficients. So,  
 114 the optimization problem becomes

$$115 \quad (1.7) \quad \text{Minimize } \left\| \sum_{j=0}^{m_k} \alpha_j^k \mathbf{F}_{k-m_k+j} \right\| \text{ subject to } \sum_{j=0}^{m_k} \alpha_j^k = 1, \alpha_j^k \geq 0.$$

116 In [21] the authors present an example where EDIIS does well and both Picard and Anderson acceleration  
 117 fail and another example where Anderson acceleration is successful and EDIIS, while converging, does not  
 118 perform as well. We present another such example in § 3. One reason why EDIIS might perform worse than  
 119 Anderson acceleration could be that the optimization problem (1.7) for EDIIS has a more restricted feasible  
 120 set and therefore a larger optimal value.

121 **2. Convergence Results.** Our global convergence is Theorem 2.1. The proof does not require differ-  
 122 entiability, but the convergence speed estimate is very pessimistic with an r-factor of  $c^{1/(m+1)}$ . We follow  
 123 the global theorem with a local theorem that shows how the convergence behavior becomes locally r-linear  
 124 with r-factor  $c$ , improving on the local results in [41].

### 125 2.1. Global Convergence.

126 **THEOREM 2.1.** *Let  $\mathbf{G}$  be a contraction on a convex set  $D \subset \mathbb{R}^N$  with contractivity constant  $c$ . Let  $\mathbf{u}^*$  be*  
 127 *the unique fixed point of  $\mathbf{G}$  in  $D$ . Then for any  $\mathbf{u}_0 \in D$ , EDIIS( $m$ ) converges to  $\mathbf{u}^*$  r-linearly with r-factor*

$$128 \quad \hat{c} = c^{1/(m+1)}.$$

129 *In fact,*

$$130 \quad (2.1) \quad \|\mathbf{e}_k\| \leq \hat{c}^k \|\mathbf{e}_0\|.$$

131 *Proof.* The proof does not use the optimality properties of the coefficients and only requires that the  
 132 iteration  $\{\mathbf{u}_k\}$  have the form

$$133 \quad (2.2) \quad \mathbf{u}_{k+1} = \sum_{j=0}^{m_k} \alpha_j^k \mathbf{G}(\mathbf{u}_{k-m_k+j}),$$

134 where  $m_k \leq m$ ,  $\alpha_j^k \geq 0$ , and  $\sum_{j=0}^{m_k} \alpha_j^k = 1$ .

135 We induct on  $k$ . Clearly (2.1) holds for both  $m_k = 0$ , by definition, and  $k = 1, m_k = 0$  because the  
 136 iteration in that case is a single Picard iteration (*i. e.* one step of Anderson(0)). Assume that the result  
 137 holds for  $k \leq K$ . Then (2.2) and  $\sum_{j=0}^{m_K} \alpha_j^K = 1$  imply that

$$138 \quad \mathbf{e}_{K+1} = \sum_{j=0}^{m_K} \alpha_j^K (\mathbf{G}(\mathbf{u}_{K-m_K+j}) - \mathbf{u}^*).$$

139 Note that since  $\alpha_j^K \geq 0$ ,  $\sum_{j=0}^{m_K} \alpha_j^K = 1$ ,  $\hat{c} < 1$ , and  $m_K \leq m$ , we have

$$140 \quad \sum_{j=0}^{m_K} \alpha_j^K \hat{c}^{K-m_K+j} \leq \hat{c}^{K-m}.$$

141 Hence

$$\begin{aligned} \|\mathbf{e}_{K+1}\| &\leq \sum_{j=0}^{m_K} \alpha_j^K \|\mathbf{G}(\mathbf{u}_{K-m_K+j}) - \mathbf{u}^*\| \\ &\leq \sum_{j=0}^{m_K} \alpha_j^K c \|\mathbf{u}_{K-m_K+j} - \mathbf{u}^*\| \\ &\leq c \sum_{j=0}^{m_K} \alpha_j^K \hat{c}^{K-m_K+j} \|\mathbf{e}_0\| \leq c \hat{c}^{K-m} \|\mathbf{e}_0\| \leq \hat{c}^{K+1} (c \hat{c}^{-m-1}) \|\mathbf{e}_0\| = \hat{c}^{K+1} \|\mathbf{e}_0\|. \end{aligned} \quad \square$$

142

143 Theorem 2.1 implies that for any  $\delta > 0$  there is  $K$  such that all iterations  $\{\mathbf{u}_k\}_{k \geq K}$  are in the set

$$144 \quad \mathcal{B}(\delta) = \{\mathbf{u} \mid \|\mathbf{u} - \mathbf{u}^*\| \leq \delta\}.$$

145 Hence, starting an Anderson acceleration iteration after sufficiently many EDIIS iterations will result in  
 146 local convergence at the rate predicted by Theorem 1.1, which is better than (2.1) since  $\hat{r}$  can be arbitrarily  
 147 near  $c$  and does not depend on  $m$ . However, it is not clear how to decide when to restart. The main result  
 148 in § 2.2, Theorem 2.2, applies to both EDIIS and Anderson acceleration, generalizes the local convergence  
 149 result from [41] (Theorem 1.1), and says that one can simply continue with the EDIIS iteration and the local  
 150 convergence estimate for Anderson acceleration will hold.

151 **2.2. Local Convergence.** Theorem 2.2 is the local convergence result. The theorem generalizes the  
 152 result in [41] by both weakening the assumptions and improving the r-factor.

153 We will assume that an iteration begins with a history that lies in  $\mathcal{B}(\delta)$  for  $\delta$  sufficiently small. This  
 154 history could be either from the EDIIS iteration or from the Anderson acceleration iteration itself. Hence  
 155 the assumption covers not only EDIIS but also allows us to improve the convergence theory from [41]. We  
 156 will show that the residuals converge r-linearly to zero with an r-factor of  $c$ . Formally our assumption is

157 **ASSUMPTION 2.1.**  $\mathbf{G}$  is a continuously differentiable contraction on  $D \subset \mathbb{R}^N$  with contractivity constant  
 158  $c$  and  $\mathbf{u}^*$  is the unique fixed point of  $\mathbf{G}$  in  $D$ .

159 The iteration begins with  $\{\mathbf{u}_l\}_{l=0}^m \subset \mathcal{B}(\delta) \subset D$ . There are real  $\{\alpha_j^k\}_{j=0}^{m_k}$  with  $0 \leq m_k \leq \min(m, k)$  such  
 160 that

$$161 \quad \sum_{j=0}^{m_k} \alpha_j^k = 1,$$

$$162 \quad (2.3) \quad \mathbf{u}_{k+1} = \sum_{j=0}^{m_k} \alpha_j^k \mathbf{G}(\mathbf{u}_{k-m_k+j}),$$

164 and

$$165 \quad (2.4) \quad \left\| \sum_{j=0}^{m_k} \alpha_j^k \mathbf{F}(\mathbf{u}_{k-m_k+j}) \right\| \leq \|\mathbf{F}(\mathbf{u}_k)\|.$$

166 Finally, there is  $\hat{c} \in (c, 1)$  so that

$$167 \quad (2.5) \quad \|\mathbf{F}(\mathbf{u}_l)\| \leq \hat{c}^l \|\mathbf{F}(\mathbf{u}_0)\|, \text{ for } 0 \leq l \leq m.$$

168 Theorem 2.1 implies that Assumption 2.1 will hold after sufficiently many EDIIS iterations. In the  
 169 theorem there is no history if  $m = 0$  and in that case the iteration is Picard iteration. While we are  
 170 motivated by a local iteration from the EDIIS algorithm, the local theory does not require that the coefficients  
 171 be nonnegative.

172 Assumption 2.1 weakens the ones in [41] in two ways. The first is that we no longer assume that  $\mathbf{G}$  is  
 173 Lipschitz continuously differentiable. The second is that we do not assume that the coefficients  $\{\alpha_j^k\}$  come  
 174 from any particular optimization problem, only that the linear combination of residuals has norm no larger  
 175 than that of the most recent residual.

176 The idea of the analysis is that as the iteration converges, the upper bound for the r-factor will approach  
 177  $c$  and therefore the r-factor is no larger than  $c$ . In the case where there is no history, this fact was implicit in  
 178 the results from [41]. Adding the history makes the bookkeeping more difficult and the proof of Theorem 2.2  
 179 must account for that.

180 **THEOREM 2.2.** Let Assumption 2.1 hold. Assume that there is  $M_\alpha$  such that

$$181 \quad (2.6) \quad \sum_{j=0}^{m_k} |\alpha_j^k| \leq M_\alpha,$$

182 for all  $k \geq 0$ . Then if  $\delta$  is sufficiently small, the iteration given by (2.3) and (2.4) converges to  $\mathbf{u}^*$  and

$$183 \quad (2.7) \quad \limsup_{k \rightarrow \infty} \left( \frac{\|\mathbf{F}(\mathbf{u}_k)\|}{\|\mathbf{F}(\mathbf{u}_0)\|} \right)^{1/k} \leq c.$$

184 *Proof.* Let  $0 < \epsilon < \hat{c} - c$ . We will show that for  $\|\mathbf{e}_0\|$  sufficiently small,

$$185 \quad (2.8) \quad \limsup_{k \rightarrow \infty} \left( \frac{\|\mathbf{F}(\mathbf{u}_k)\|}{\|\mathbf{F}(\mathbf{u}_0)\|} \right)^{1/k} \leq c + \epsilon.$$

186 This will complete the proof since  $\epsilon$  is arbitrary and we can restart the proof once we have  $m$  vectors in the  
187 history which are near enough to  $\mathbf{u}^*$  to reduce  $\epsilon$  further.

188 We induct on  $k$ . Define  $L = (c/\hat{c})^m$ . We will show that

$$189 \quad (2.9) \quad \|\mathbf{F}(\mathbf{u}_k)\| \leq L(c + \epsilon)^k \|\mathbf{F}(\mathbf{u}_0)\|,$$

190 for all  $k$ . Our assumption on the history that  $\|\mathbf{F}(\mathbf{u}_l)\| \leq \hat{c}^l \|\mathbf{F}(\mathbf{u}_0)\|$  implies that (2.9) holds for  $0 \leq k \leq m$ .  
191 Now suppose that (2.9) holds for all  $0 \leq l \leq k$  with  $k \geq m$ .

192 We will establish the bound for  $k + 1$ . The analysis has three steps. We first set  $\delta$  small enough for the  
193 iteration to remain in  $D$ . We then derive an estimate for  $\mathbf{F}(\mathbf{u}_{k+1})$  and finally use that estimate to continue  
194 the induction.

195 **Step 1, initialization of  $\delta$ :** Since  $\mathbf{G}'$  is continuous in  $D$ , there is a nondecreasing function  $\rho \in C[0, \infty)$   
196 with  $\rho(0) = 0$  so that

$$197 \quad (2.10) \quad \|\mathbf{G}'(\mathbf{u}) - \mathbf{G}'(\mathbf{u}^*)\| \leq \rho(\|\mathbf{e}\|)$$

198 for all  $\mathbf{u} \in D$ . This implies that for all  $\mathbf{u} \in D$ ,

$$199 \quad (2.11) \quad \mathbf{G}(\mathbf{u}) = \mathbf{G}(\mathbf{u}^*) + \int_0^1 \mathbf{G}'(\mathbf{u}^* + t\mathbf{e})\mathbf{e} dt = \mathbf{u}^* + \mathbf{G}'(\mathbf{u}^*)\mathbf{e} + \Delta(\mathbf{e}),$$

200 where

$$201 \quad \|\Delta(\mathbf{e})\| \leq \rho(\|\mathbf{e}\|)\|\mathbf{e}\|.$$

202 Contractivity of  $\mathbf{G}$  implies that

$$203 \quad \|\mathbf{F}(\mathbf{u})\|/(1 + c) \leq \|\mathbf{e}\| \leq \|\mathbf{F}(\mathbf{u})\|/(1 - c).$$

204 Assumption 2.1 implies that

$$205 \quad \mathcal{B}(\delta) \cap \{\mathbf{u} \mid \|\mathbf{F}(\mathbf{u})\| \leq \|\mathbf{F}(\mathbf{u}_0)\|\} \subset D.$$

206 Reduce  $\delta$  if necessary so that

$$207 \quad (2.12) \quad \rho \left( M_\alpha L(c + \epsilon)^{k-m} \delta \frac{1+c}{1-c} \right) \leq \frac{c^{m+1}(1-c)}{2M_\alpha} \left( 1 - \frac{c}{c+\epsilon} \right).$$

208 This implies that

$$209 \quad (2.13) \quad \mathbf{w}_k = \sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j} \in D$$

210 for sufficiently small  $\delta$  because

$$211 \quad (2.14) \quad \begin{aligned} \|\mathbf{w}_k - \mathbf{u}^*\| &\leq \sum_{j=0}^{m_k} |\alpha_j^k| \|\mathbf{e}_{k-m_k+j}\| \\ &\leq M_\alpha L(c + \epsilon)^{k-m} \|\mathbf{F}(\mathbf{u}_0)\|/(1 - c) \leq M_\alpha L(c + \epsilon)^{k-m} \delta(1 + c)/(1 - c). \end{aligned}$$

212 **Step 2, estimation of  $\mathbf{F}(\mathbf{u}_{k+1})$ :** We may write for  $k \geq m - 1$ ,

$$\begin{aligned}
 213 \quad \mathbf{F}(\mathbf{u}_{k+1}) &= \mathbf{G}(\mathbf{u}_{k+1}) - \mathbf{u}_{k+1} \\
 &= \mathbf{G}(\mathbf{u}_{k+1}) - \mathbf{G}(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}) + \mathbf{G}(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}) - \mathbf{u}_{k+1}.
 \end{aligned}$$

214 We will estimate the two parts of the sum

$$215 \quad A_k = \mathbf{G}(\mathbf{u}_{k+1}) - \mathbf{G}\left(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}\right)$$

216 and

$$217 \quad B_k = \mathbf{G}\left(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}\right) - \mathbf{u}_{k+1}$$

218 separately.

219 Using only contractivity of  $\mathbf{G}$  and (2.4) we have

$$\begin{aligned}
 \|A_k\| &= \|\mathbf{G}(\mathbf{u}_{k+1}) - \mathbf{G}(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j})\| \\
 &\leq c \|\mathbf{u}_{k+1} - \sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}\| \\
 220 \quad (2.15) \quad &= c \|\sum_{j=0}^{m_k} \alpha_j^k (\mathbf{G}(\mathbf{u}_{k-m_k+j}) - \mathbf{u}_{k-m_k+j})\| \\
 &= c \|\sum_{j=0}^{m_k} \alpha_j^k \mathbf{F}(\mathbf{u}_{k-m_k+j})\| \leq c \|\mathbf{F}(\mathbf{u}_k)\|.
 \end{aligned}$$

221 We now estimate  $B_k$ . Using (2.12) we have for all  $\mathbf{u} \in D$  with

$$222 \quad \|\mathbf{e}\| \leq M_\alpha L(c + \epsilon)^{k-m} \delta(1 + c)/(1 - c)$$

223

$$\begin{aligned}
 \|\Delta(\mathbf{e})\| &\leq \rho(\|\mathbf{e}\|) \|\mathbf{F}(\mathbf{u})\|/(1 - c) \\
 224 \quad (2.16) \quad &\leq \rho(M_\alpha L(c + \epsilon)^{k-m} \delta(1 + c)/(1 - c)) \|\mathbf{F}(\mathbf{u})\|/(1 - c) \\
 &\leq \frac{c^{m+1}}{2M_\alpha} \left(1 - \frac{c}{c+\epsilon}\right) \|\mathbf{F}(\mathbf{u})\|.
 \end{aligned}$$

225 The final stage in the proof is to show that, reducing  $\delta$  if needed,

$$226 \quad (2.17) \quad \|B_k\| \leq L(c + \epsilon)^{k+1} \left(1 - \frac{c}{c + \epsilon}\right) \|\mathbf{F}(\mathbf{u}_0)\|.$$

227 Recall that

$$\begin{aligned}
 228 \quad B_k &= \mathbf{G}(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}) - \mathbf{u}_{k+1} \\
 &= \mathbf{G}(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}) - \sum_{j=0}^{m_k} \alpha_j^k \mathbf{G}(\mathbf{u}_{k-m_k+j}).
 \end{aligned}$$

229 We use (2.11) to obtain

$$\begin{aligned}
 \mathbf{G}(\sum_{j=0}^{m_k} \alpha_j^k \mathbf{u}_{k-m_k+j}) &= \mathbf{G}(\mathbf{w}_k) = \mathbf{u}^* + \mathbf{G}'(\mathbf{u}^*) \sum_{j=0}^{m_k} \alpha_j^k \mathbf{e}_{k-m_k+j} + \Delta(\mathbf{w}_k - \mathbf{u}^*) \\
 230 \quad &= \sum_{j=0}^{m_k} \alpha_j^k (\mathbf{u}^* + \mathbf{G}'(\mathbf{u}^*) \mathbf{e}_{k-m_k+j}) + \Delta(\mathbf{w}_k - \mathbf{u}^*) \\
 &= \sum_{j=0}^{m_k} \alpha_j^k \mathbf{G}(\mathbf{u}_{k-m_k+j}) + \sum_{j=0}^{m_k} \alpha_j^k \Delta(\mathbf{e}_{k-m_k+j}) + \Delta(\mathbf{w}_k - \mathbf{u}^*).
 \end{aligned}$$

231 Hence

$$232 \quad \|B_k\| \leq \sum_{j=0}^{m_k} |\alpha_j^k| \|\Delta(\mathbf{e}_{k-m+1})\| + \|\Delta(\mathbf{w}_k - \mathbf{u}^*)\|.$$

233 We will estimate terms separately. First

$$\begin{aligned} \sum_{j=0}^{m_k} |\alpha_j^k| \|\Delta(\mathbf{e}_{k-m+1})\| &\leq \frac{c^{m+1}}{2M_\alpha} \left(1 - \frac{c}{c+\epsilon}\right) \sum_{j=0}^{m_k} |\alpha_j^k| \|\mathbf{F}(\mathbf{u}_{k-m_k+j})\| \\ &\leq \frac{c^{m+1}}{2M_\alpha} \left(1 - \frac{c}{c+\epsilon}\right) \sum_{j=0}^{m_k} |\alpha_j^k| L(c+\epsilon)^{k-m_k+j} \|\mathbf{F}(\mathbf{u}_0)\| \\ 234 \quad (2.18) \quad &\leq \frac{c^{m+1}}{2} \left(1 - \frac{c}{c+\epsilon}\right) L(c+\epsilon)^{k-m_k} \|\mathbf{F}(\mathbf{u}_0)\| \\ &\leq (L/2)(c+\epsilon)^{k+1} \left(1 - \frac{c}{c+\epsilon}\right) \|\mathbf{F}(\mathbf{u}_0)\|. \end{aligned}$$

235 Finally, using (2.14) and (2.16),

$$\begin{aligned} \|\Delta(\mathbf{w}_k - \mathbf{u}^*)\| &\leq \rho(\|\mathbf{w}_k - \mathbf{u}^*\|) \|\mathbf{F}(\mathbf{w}_k)\| / (1-c) \\ 236 \quad (2.19) \quad &\leq \rho(M_\alpha L(c+\epsilon)^{k-m} \delta(1+c) / (1-c)) M_\alpha L(c+\epsilon)^{k-m} \|\mathbf{F}(\mathbf{u}_0)\| / (1-c) \\ &\leq (L/2)(c+\epsilon)^{k+1} \left(1 - \frac{c}{c+\epsilon}\right) \|\mathbf{F}(\mathbf{u}_0)\|. \end{aligned}$$

237 Adding the two estimates (2.18) and (2.19) leads to (2.17).

238 **Step 3, continuation of the induction:** Combining (2.15), (2.17), (2.9), and the induction hypothe-  
239 ses, we have

$$\begin{aligned} \|\mathbf{F}(\mathbf{u}_{k+1})\| &\leq c \|\mathbf{F}(\mathbf{u}_k)\| + L(c+\epsilon)^{k+1} \left(1 - \frac{c}{c+\epsilon}\right) \|\mathbf{F}(\mathbf{u}_0)\| \\ 240 \quad (2.20) \quad &\leq \left( Lc(c+\epsilon)^k + L(c+\epsilon)^{k+1} \left(1 - \frac{c}{c+\epsilon}\right) \right) \|\mathbf{F}(\mathbf{u}_0)\| \\ &\leq L(c+\epsilon)^{k+1} \|\mathbf{F}(\mathbf{u}_0)\|. \end{aligned}$$

241 This implies (2.8), which in turn implies (2.7) because  $\epsilon$  is arbitrary.  $\square$

242 Theorem 2.2 and nonsingularity of  $\mathbf{F}'(\mathbf{u}^*)$  also imply r-linear convergence of the errors with r-factor  $c$ .  
243 This extends and sharpens (1.6).

244 **COROLLARY 2.3.** *Let the assumptions of Theorem 2.2 hold. If  $\mathbf{F}'(\mathbf{u}^*)$  is nonsingular then*

$$245 \quad (2.21) \quad \limsup_{k \rightarrow \infty} \left( \frac{\|\mathbf{e}_k\|}{\|\mathbf{e}_0\|} \right)^{1/k} \leq c.$$

246 *Proof.* We will use Lemma 5.2.1 from [19], which states that if  $\mathbf{u}$  is sufficiently near  $\mathbf{u}^*$  and  $\mathbf{F}'(\mathbf{u}^*)$  is  
247 nonsingular, then

$$248 \quad \frac{\|\mathbf{e}\|}{\|\mathbf{e}_0\|} \leq 4 \|\mathbf{F}'(\mathbf{u}^*)\| \|\mathbf{F}'(\mathbf{u}^*)^{-1}\| \frac{\|\mathbf{F}(\mathbf{u})\|}{\|\mathbf{F}(\mathbf{u}_0)\|}.$$

249 Hence

$$250 \quad \limsup_{k \rightarrow \infty} \left( \frac{\|\mathbf{e}_k\|}{\|\mathbf{e}_0\|} \right)^{1/k} \leq \lim_{k \rightarrow \infty} \left( 4 \|\mathbf{F}'(\mathbf{u}^*)\| \|\mathbf{F}'(\mathbf{u}^*)^{-1}\| \right)^{1/k} \limsup_{k \rightarrow \infty} \left( \frac{\|\mathbf{F}(\mathbf{u}_k)\|}{\|\mathbf{F}(\mathbf{u}_0)\|} \right)^{1/k} \leq c,$$

251 which is (2.21).  $\square$

**3. Numerical Example.** We will use an example [41] to show how the actual performance of EDIIS and Anderson acceleration can differ, even though the theoretical limiting convergence estimates are identical. Another point of this section is that the solver for the optimization problem can significantly affect the results.

The results in [21] also illustrate this point. Our example is simple enough to directly compare the iteration histories for Picard iteration, EDIIS, and Anderson with the worst-case prediction given by the contractivity constant. We find that when Anderson acceleration performs well, as it does in this example, EDIIS offers no advantage. Moreover, the additional constraint on the optimization problem for the coefficients leads to slower convergence, exactly matching Picard iteration in this case.

The optimization problem for EDIIS requires more care than the linear least squares problem one must solve for Anderson acceleration. The reason for this is that one cannot simply use a QR factorization to solve (1.4). Instead one must apply a more sophisticated iterative solver. The approach of [21] is a direct examination of the boundary of the feasible simplex, which is not practical for a depth much greater than  $m = 3$ . Since  $m$  is small in practice, expressing the optimization problem as a bound-constrained quadratic program is an efficient alternative. [26, 27] survey the literature on this topic. For example a bound-constrained quadratic programming code such as the MINQ [29] code is a reasonable choice. However this approach squares the condition number and can (and did in our testing) result in a singular or nearly singular KKT system and failure of the optimization code’s internal linear solvers. The method of [6], while still squaring the condition number, is more robust and terminated without error for this example. The classic method from [14] uses an active set method and the QR factorization to avoid using the normal equations. The approach in [14] performed better in the example here, where the least squares coefficient matrix for the optimization problem is ill-conditioned [41].

The example is the midpoint rule discretization of the Chandrasekhar H-equation [3, 5].

$$(3.1) \quad \mathcal{F}(H)(\mu) = H(\mu) - \left(1 - \frac{\omega}{2} \int_0^1 \frac{\mu H(\nu) d\nu}{\mu + \nu}\right)^{-1} = 0.$$

We seek a solution  $H^* \in C[0, 1]$ . When the parameter  $\omega$  is important we will write  $H^*$  as a function  $H^*(\mu, \omega)$  of both  $\mu$  and  $\omega$ .

The integral equation and its midpoint discretization share the properties that the fixed point map

$$\mathcal{G}(H)(\mu) = \left(1 - \frac{\omega}{2} \int_0^1 \frac{\mu H(\nu) d\nu}{\mu + \nu}\right)^{-1}$$

is a contraction for  $0 \leq \omega < 1$ , but not for  $\omega = 1$ . The Fréchet derivative (and the Jacobian for the discrete case) is singular at the solution for  $\omega = 1$ , which is a simple fold singularity [17, 28].

In this section we will compare the performance of Picard iteration, Anderson acceleration, and EDIIS for the case  $\omega = .5$  on an  $N = 100$  point mesh. We terminated the iteration when the residual had decreased by a factor of  $10^{-12}$ .

One interesting result from [41] is that Anderson(m) is more efficient than Newton’s method for this example, even in the singular case. In the context of this paper it is important to note that Picard iteration converges faster than one would expect from estimating the contractivity parameter by the spectral radius of the Fréchet derivative of  $\mathcal{G}$  at the solution, which is a lower bound for the operator norm of  $\mathcal{G}$ . From [41]

$$\rho(\mathcal{G}'(H^*)) = 1 - \sqrt{1 - \omega} \approx .293.$$

However [2, 18, 20], the solution is analytic in  $\omega$  and Picard iteration exploits that property to obtain q-linear convergence with q-factor  $\leq \rho(\mathcal{G}'(H^*))$  and much less for small  $\omega$ . In fact, if

$$H^*(\mu, \omega) = \sum_{m=0}^{\infty} \omega^m a_m(\mu)$$

is the Taylor expansion of  $H^*$  in  $\omega$  then the coefficient functions  $\{a_m(\mu)\}$  are nonnegative for  $0 \leq \mu \leq 1$ . Moreover the series converges for  $\omega = 1$ . Hence, if  $H_k$  is the  $k$ th Picard iteration and  $H_0 \equiv 0$ , then for all  $k \geq 0$  and  $\omega, \mu \in [0, 1]$ ,

$$H_k(\mu, \omega) \leq H_{k+1}(\mu, \omega) \leq H^*(\mu, \omega).$$

296 All of the above statements about the singularity at  $\omega = 1$ , the spectral radius of the Fréchet derivative,  
 297 and the performance of Picard iteration apply to the discrete problem

$$298 \quad (3.2) \quad \mathbf{G}(\mathbf{h})_i = \left( 1 - \frac{\omega}{2N} \sum_{j=1}^N \frac{h_j \mu_i}{\mu_i + \mu_j} \right)^{-1}, \quad 1 \leq i \leq N.$$

299 In (3.2)  $\mu_i = (i - 1/2)/N$  is the  $i$ th quadrature node for the  $N$  point composite midpoint rule, the vector  
 300  $\mathbf{h}^*$  is the solution of the discrete problem  $\mathbf{h} = \mathbf{G}(\mathbf{h})$ ,  $\mathbf{G}(\mathbf{h}^*)_i$  is the  $i$ th component of  $\mathbf{G}(\mathbf{h}^*)$ , and the  $i$ th  
 301 component of  $\mathbf{h}^*$  is  $h_i^* \approx H^*(\mu_i)$ .

302 As noted above, the optimization problem (1.7) for EDIIS is harder than the one for Anderson accelera-  
 303 tion and the choice of solver can be important. We compare the method of [14], as implemented in the Matlab  
 304 `lsqlin` code with the ‘active-set’ option, with the method from [6], as implemented with ‘interior-point’ op-  
 305 tion in `lsqlin`. The method of [6] uses the normal equations and did exhibit problems with ill-conditioning.  
 306 The computations were done on an Apple Macintosh running MAC OS 10.13.6 with Matlab 2017a. The  
 307 ‘active-set’ option was removed with Matlab 2017b. The codes that generated Table 3.1 and Figure 3.1 are  
 308 supplementary materials for this paper.

309 In the left plot of Figure 3.1 we compare Picard iteration, Anderson acceleration, and EDIIS with the  
 310 active-set option (EDIIS-A) and the interior-point option (EDIIS-I). The depth was  $m = 3$  for the Anderson  
 311 and EDIIS computations. Picard iteration and EDIIS-A are identical. The optimization problem for EDIIS  
 312 cannot match the results from Anderson acceleration, which has fairly large negative coefficients. Rather,  
 313 EDIIS-A finds that the coefficients for Picard iteration are optimal.

314 Table 3.1 compares  $\rho(\mathcal{G}'(H^*))$  to the r-factors of the residuals for Anderson acceleration, Picard iteration,  
 315 and EDIIS. We estimate the r-factors by

$$316 \quad \left( \frac{\|\mathbf{F}(\mathbf{h}_k)\|}{\|\mathbf{F}(\mathbf{h}_0)\|} \right)^{1/k}$$

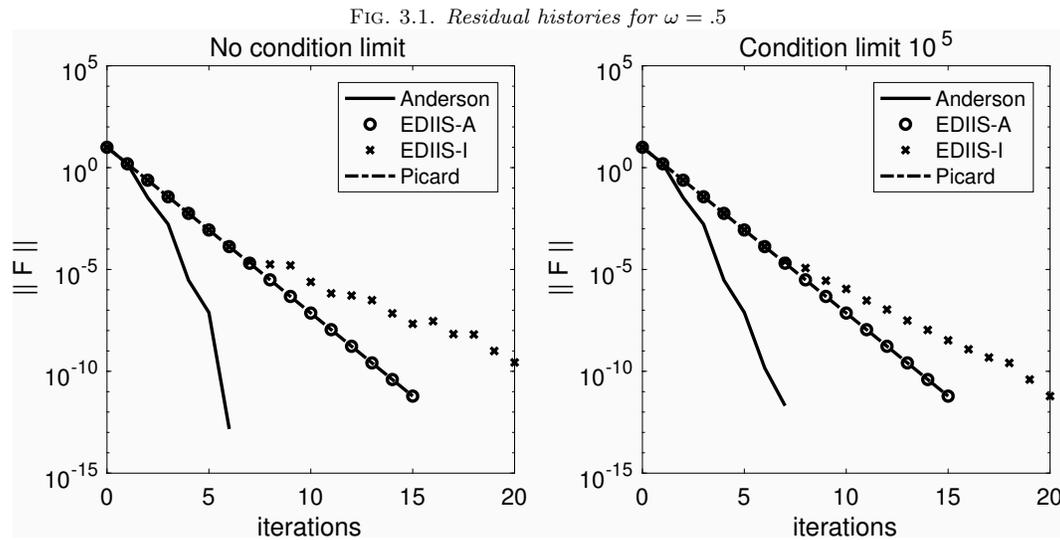
317 where the final iteration upon termination is  $\mathbf{h}_k$ . Note that, as discussed above, the q-factor for Picard  
 318 iteration is smaller than the spectral radius. Anderson acceleration also does better than the theory predicts  
 319 and, in fact, is more efficient than Newton-GMRES [41].

320 EDIIS-I is the only one of the methods which is sensitive to the ill-conditioning of the optimization  
 321 problem. We examined this sensitivity by solving the problem twice, once with no limit on the condition  
 322 number and again by reducing  $m_k$  if necessary to limit the condition number to  $10^5$ . This has no effect  
 323 on EDIIS-A and slightly slows Anderson acceleration down. We show the residual histories in Figure 3.1,  
 324 where one can clearly see the effect of limiting the condition number. As reported in [41], the optimization  
 325 problem becomes more ill-conditioned as the iteration progresses. The figures show that the convergence of  
 326 EDIIS-I degrades at the 6th iteration, but to a lesser degree when the condition number is limited. Note  
 327 that the estimated r-factor seems to stabilize near the end and is, in the condition number limited case, back  
 328 to Picard iteration for the final three iterations, albeit from a worse starting point.

TABLE 3.1  
 Convergence r-Factors

Anderson	Picard	EDIIS-A	EDIIS-I	$\rho(\mathcal{G}'(H^*))$
No condition limit				
1.06e-02	1.72e-01	1.72e-01	2.62e-01	2.93e-01
Condition limit $10^5$				
2.59e-02	1.72e-01	1.72e-01	2.62e-01	2.93e-01

329 **4. Conclusions.** The EDIIS algorithm was designed to improve the global convergence properties of  
 330 the DIIS algorithm, which is also known as Anderson acceleration. We prove global convergence of the  
 331 iteration and prove a local convergence result that applies to both EDIIS and Anderson acceleration and  
 332 improves the results in [41]. We observe, as did the inventors of the method [21], that the unmodified version  
 333 of Anderson acceleration can have better local convergence in practice.



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337

## REFERENCES

- 338 [1] D. G. ANDERSON, *Iterative Procedures for Nonlinear Integral Equations*, Journal of the ACM, 12 (1965), pp. 547–560,  
 339 doi:10.1145/321296.321305, <http://portal.acm.org/citation.cfm?doid=321296.321305>.
- 340 [2] P. B. BOSMA AND W. A. DEROOIJ, *Efficient methods to calculate Chandrasekhar’s H-functions*, Astron. Astrophys., 126  
 341 (1983), p. 283.
- 342 [3] I. W. BUSBRIDGE, *The Mathematics of Radiative Transfer*, no. 50 in Cambridge Tracts, Cambridge Univ. Press, Cam-  
 343 bridge, 1960.
- 344 [4] N. N. CARLSON AND K. MILLER, *Design and application of a gradient weighted moving finite element code I: In one*  
 345 *dimension*, SIAM J. Sci. Comput., 19 No. 3 (1998), pp. 766–798.
- 346 [5] S. CHANDRASEKHAR, *Radiative Transfer*, Dover, New York, 1960.
- 347 [6] T. F. COLEMAN AND Y. LI, *On the convergence of interior-reflective Newton methods for nonlinear minimization subject*  
 348 *to bounds*, Math. Prog., 67 (1994), pp. 189–224.
- 349 [7] A. M. COLLIER, A. C. HINDMARSH, R. SERBAN, AND C. S. WOODWARD, *User documentation for KINSOL v2.8.0*, Tech.  
 350 Report UCRL-SM-208116, Lawrence Livermore National Laboratory, 2015.
- 351 [8] CVX RESEARCH, INC., *CVX: Matlab software for disciplined convex programming, version 2.0*. <http://cvxr.com/cvx>,  
 352 Aug. 2012.
- 353 [9] P. H. DEDRICHS AND R. ZELLER, *Self-consistency iterations in electronic-structure calculations*, Phys. Rev. B, 28 (1983),  
 354 pp. 5462–5472.
- 355 [10] J. DEGROOTE, K.-J. BATHE, AND J. VIERENDEELS, *Performance of a new partitioned procedure versus a monolithic*  
 356 *procedure in fluid-structure interaction*, Computers and Structures, 97 (2009), pp. 793–801.
- 357 [11] H.-R. FANG AND Y. SAAD, *Two classes of multisecond methods for nonlinear acceleration*, Numerical Linear Algebra with  
 358 Applications, 16 (2009), pp. 197–221, doi:10.1002/nla.
- 359 [12] M. J. FRISCH, G. W. TRUCKS, H. B. SCHLEGEL, G. E. SCUSERIA, M. A. ROBB, J. R. CHEESEMAN, G. SCALMANI,  
 360 V. BARONE, B. MENNUCCI, G. A. PETERSSON, H. NAKATSUJI, M. CARICATO, X. LI, H. P. HRATCHIAN, A. F. IZ-  
 361 MAYLOV, J. BLOINO, G. ZHENG, J. L. SONNENBERG, M. HADA, M. EHARA, K. TOYOTA, R. FUKUDA, J. HASEGAWA,  
 362 M. ISHIDA, T. NAKAJIMA, Y. HONDA, O. KITAO, H. NAKAI, T. VREVEN, J. A. MONTGOMERY JR., J. E. PERALTA,  
 363 F. OGLIARO, M. BEARPARK, J. J. HEYD, E. BROTHERS, K. N. KUDIN, V. N. STAROVEROV, R. KOBAYASHI, J. NOR-  
 364 MAND, K. RAGHAVACHARI, A. RENDELL, J. C. BURANT, S. S. IYENGAR, J. TOMASI, M. COSSI, N. REGA, J. M.  
 365 MILLAM, M. KLENE, J. E. KNOX, J. B. CROSS, V. BAKKEN, C. ADAMO, J. JARAMILLO, R. GOMPERS, R. E. STRAT-  
 366 MANN, O. YAZYEV, A. J. AUSTIN, R. CAMMI, C. POMELLI, J. W. OCHTERSKI, R. L. MARTIN, K. MOROKUMA, V. G.  
 367 ZAKRZEWSKI, G. A. VOTH, P. SALVADOR, J. J. DANNENBERG, S. DAPPRICH, A. D. DANIELS, Ö. FARKAS, J. B.  
 368 FORESMAN, J. V. ORTIZ, J. CIOSLOWSKI, AND D. J. FOX, *Gaussian 09, Revision A.1*, 2009.
- 369 [13] D. J. GARDNER, C. S. WOODWARD, D. R. REYNOLDS, G. HOMMES, S. AUBREY, AND A. ARSNELIS, *Implicit integration*  
 370 *methods for dislocation dynamics*, Modelling Simul. Mater. Sci. Eng., 23 (2015), p. 025006 (31pp).
- 371 [14] G. H. GOLUB AND M. A. SAUNDERS, *Linear least squares and quadratic programming*, Tech. Report CS 134, Stanford  
 372 University, 1969.

- [15] R. HAELTERMAN, J. DEGROOTE, D. VAN HEULE, AND J. VIERENDEELS, *The quasi-newton least squares method: A new and fast secant method analyzed for linear systems*, SIAM J. Numer. Anal., 47 (2009), pp. 2347–2368.
- [16] S. HAMILTON, M. BERRILL, K. CLARNO, R. PAWLOWSKI, A. TOTH, C. T. KELLEY, T. EVANS, AND B. PHILIP, *An assessment of coupling algorithms for nuclear reactor core physics simulations*, Journal of Computational Physics, 311 (2016), pp. 241–257.
- [17] H. B. KELLER, *Lectures on Numerical Methods in Bifurcation Theory*, Tata Institute of Fundamental Research, Lectures on Mathematics and Physics, Springer-Verlag, New York, 1987.
- [18] C. T. KELLEY, *Solution of  $H$ -equations by iteration*, SIAM J. Math. Anal., 10 (1979), pp. 844–849.
- [19] C. T. KELLEY, *Iterative Methods for Linear and Nonlinear Equations*, no. 16 in Frontiers in Applied Mathematics, SIAM, Philadelphia, 1995.
- [20] C. T. KELLEY AND T. W. MULLIKIN, *Solution by iteration of  $H$ -equations in multigroup neutron transport*, J. Math. Phys., 19 (1978), pp. 500–501.
- [21] K. N. KUDIN, G. E. SCUSERIA, AND E. CANCÈS, *A black-box self-consistent field convergence algorithm: One step closer*, Journal of Chemical Physics, 116 (2002), pp. 8255–8261, doi:10.1063/1.1470195.
- [22] L. LIN AND C. YANG, *Elliptic preconditioner for accelerating the self-consistent field iteration in Kohn-Sham density functional theory*, SIAM J. Sci. Comp., 35 (2013), pp. S277–S298.
- [23] F. LINDNER, M. MEHL, K. SCHEUFELE, AND B. UEKERMANN, *A comparison of various quasi-newton schemes for partitioned fluid-structure interaction*, in ECCOMAS Coupled Problems in Science and Engineering, Venice, B. A. Schrefler, E. Oñate, and M. Papadarakis, eds., Barcelona, 2015, DIMNE, pp. 477–488.
- [24] P. A. LOTT, H. F. WALKER, C. S. WOODWARD, AND U. M. YANG, *An accelerated Picard method for nonlinear systems related to variably saturated flow*, Advances in Water Resources, 38 (2012), pp. 92–101.
- [25] K. MILLER, *Nonlinear Krylov and moving nodes in the method of lines*, Journal of Computational and Applied Mathematics, 183 (2005), pp. 275–287.
- [26] J. J. MORÉ AND G. TORALDO, *Algorithms for bound constrained quadratic programming problems*, Numer. Math., 55 (1989), pp. 377–400.
- [27] J. J. MORÉ AND G. TORALDO, *On the solution of large quadratic programming problems with bound constraints*, SIAM J. Optim., 1 (1991), pp. 93–113.
- [28] T. W. MULLIKIN, *Some probability distributions for neutron transport in a half space*, J. Appl. Prob., 5 (1968), pp. 357–374.
- [29] A. NEUMAIER, *MINQ - General Definite and Bound Constrained Indefinite Quadratic Programming*, 1998, <http://www.mat.univie.ac.at/~neum/software/minq/>.
- [30] C. W. OOSTERLEE AND T. WASHIO, *Krylov subspace acceleration for nonlinear multigrid schemes*, SIAM J. Sci. Comput., 21 (2000), pp. 1670–1690.
- [31] F. A. POTRA AND H. ENGLER, *A characterization of the behavior of the Anderson acceleration on linear problems*, Lin. Alg. Appl., 438 (2013), pp. 1002–1011.
- [32] P. PULAY, *Convergence acceleration of iterative sequences. The case of SCF iteration*, Chemical Physics Letters, 73 (1980), pp. 393–398, <http://www.sciencedirect.com/science/article/pii/0009261480803964>.
- [33] P. PULAY, *Improved SCF convergence acceleration*, Journal of Computational Chemistry, 3 (1982), pp. 556–560, <http://onlinelibrary.wiley.com/doi/10.1002/jcc.540030413/abstract>.
- [34] T. ROHWEDDER AND R. SCHNEIDER, *An analysis for the DIIS acceleration method used in quantum chemistry calculations*, Journal of Mathematical Chemistry, 49 (2011), pp. 1889–1914, <http://www.springerlink.com/index/10.1007/s10910-011-9863-y>.
- [35] Y. SAAD, J. R. CHELIKOWSKY, AND S. M. SHONTZ, *Numerical methods for electronic structure calculations of materials*, SIAM Review, 52 (2010), pp. 3–54.
- [36] H. B. SCHLEGEL AND J. J. W. MCDUALL, *Do you have SCF stability and convergence problems?*, in Computational Advances in Organic Chemistry: Molecular Structure and Reactivity, C. Ögretir and I. G. Csizmadia, eds., Dordrecht, 1991, Kluwer, pp. 167–185.
- [37] R. SCHNEIDER, T. ROHWEDDER, A. NEELOV, AND J. BLAUERT, *Direct minimization for calculating invariant subspaces in density functional computations of the electronic structure*, Journal of Computational Mathematics, 27 (2008), pp. 360–387.
- [38] *SUNDIALS (SUite of Nonlinear and Differential/ALgebraic Solvers)*. <http://www.llnl.gov/casc/sundials>.
- [39] A. TOTH, *A Theoretical Analysis of Anderson Acceleration and Its Application in Multiphysics Simulation for Light-Water Reactors*, PhD thesis, North Carolina State University, Raleigh, North Carolina, 2016.
- [40] A. TOTH, J. A. ELLIS, T. EVANS, S. HAMILTON, C. T. KELLEY, R. PAWLOWSKI, AND S. SLATTERY, *Local improvement results for Anderson acceleration with inaccurate function evaluations*, SIAM J. Sci. Comp., 39 (2017), pp. S47–S65, doi:10.1137/16M1080677.
- [41] A. TOTH AND C. T. KELLEY, *Convergence analysis for Anderson acceleration*, SIAM J. Numer. Anal., 53 (2015), pp. 805–819, doi:10.1137/130919398.
- [42] A. TOTH, C. T. KELLEY, S. SLATTERY, S. HAMILTON, K. CLARNO, AND R. PAWLOWSKI, *Analysis of Anderson acceleration on a simplified neutronics/thermal hydraulics system*, 2015. Joint International Conference on Mathematics and Computation (M&C), Supercomputing in Nuclear Applications (SNA) and the Monte Carlo (MC) Method.
- [43] A. TOTH AND R. PAWLOWSKI, *NOX::Solvers::AndersonAcceleration Class Reference*, 2015, [https://trilinos.org/docs/dev/packages/nox/doc/html/classNOX\\_1\\_1Solver\\_1\\_1AndersonAcceleration.html](https://trilinos.org/docs/dev/packages/nox/doc/html/classNOX_1_1Solver_1_1AndersonAcceleration.html).
- [44] H. W. WALKER AND P. NI, *Anderson acceleration for fixed-point iterations*, SIAM J. Numer. Anal., 49 (2011), pp. 1715–1735.
- [45] T. WASHIO AND C. OOSTERLEE, *Krylov subspace acceleration for nonlinear multigrid schemes*, Elec. Trans. Num. Anal., 6 (1997), pp. 271–290.

- 439 [46] D. YOUNG, *Computational Chemistry: A Practical Guide for Applying Techniques to Real World Problems*, Wiley, New  
440 York, 2001.
- 441 [47] D. YOUNG, *SCF convergence and chaos theory*. <http://www.ccl.net/cca/documents/dyoung/topics-orig/converge.html>,  
442 2001.